## t d t t ndth n f n t d dSt utu n - t h u

Xiu en Zhang and Ale Zunger E, L, G, C, 80401, A(g2/F51Tf16.34070TD(B)Tj/F41Tf0.65890TD(2)Tj/F21Tf0.49720TD[())-385.8(has)-387.1(its)]TJ/F41Tf3 Li, the hR1 structure transforms to the cI16 structure at about 35 GPa, then at about 75 GPa it transforms to the C2structure, and then at about 98 GPa it transforms to the *Aba2* structure, hich persists up to at least 160 GPa; for Cs, the *Cmca* structure transforms to the *dhcp* structure at about 60 GPa, hich persists up to at least 200 GPa. The transition pressures are in agreement ith the e perimental results [3, 5]. Random search techniques [24] and the evolutionar algorithm method [25,26] ere previousl

## $SI: 7_{resp} B_{rel}^{to} = 6533 (3.2)^{rd} S_{resp}^{table} = 8240.(,I)-421)D7.63$

$$\Delta H_{\rm VD}$$
 ¼ ½HðLi,  $V_m$ Þ HðLi,  $V_1$ Þ þ½HðCs,  $V_m$ Þ

$$H\delta Cs, V_2 \flat$$
. (1)

The charge-transfer energ  $\Delta H_{\rm CT}$  is the enthalp change hen combining bcc Cs and bcc Li at the constant-volume  $V_m$  to form LiCs (B2),

$$\Delta H_{\rm CT}$$
 ¼ HðLiCs,  $V_m \flat$  HðLi,  $V_m \flat$  HðCs,  $V_m \flat$ . (2)

 $\Delta H_{\rm VD}$  and  $\Delta H_{\rm CT}$  are shon as the red line and the green line in Fig. 4(a), respectivel. We see that  $\Delta H_{\rm VD}$  is positive and increases rapidle as pressure increases, the main effect coming from the  $P \delta V_2 = V_1 b/2$  term (dashed line). At error pressure,  $\Delta H_{\rm VD}$  (b520 meV/atom) is large enough to make the Li-Cs s stem phase separating. In contrast,  $\Delta H_{\rm CT}$  is negative and decreases rapidle as pressure increases. At 160 GPa,  $\Delta H \frac{14}{4}$  258 meV/atom and  $\Delta H_{\rm CT}$ is large enough to induce ordering in Li-Cs.

The electronic structures of LiCs (*B*2) are sho n in Figs. 4(b) and 4(c). At r ero pressure the dispersions of the Li $\delta$ 1*s* $\vartheta$ , Cs $\delta$ 5*s* $\vartheta$ , and Cs $\delta$ 5*p* $\vartheta$  bands are all small. At 160 GPa, the Li $\delta$ 1*s* $\vartheta$  band still has negligible dispersion, hile the dispersion of the Cs $\delta$ 5*s* $\vartheta$ 

high-pressure compounds bet een t o elements in the same column of the periodic table (Li-Cs) is perhaps reminiscent of compound formation and charge transfer bet een Ag-K [27] or Li-Be [17], although the latter t o cases are perhaps less surprising than Li-Cs as the represent a far larger chemical disparit (hence, reactivit) bet een the constituent elements that belong to different columns in the periodic table.

Li<sub>7</sub>Cs  $(Cmmm)^{x}$ , *1D*-, *The electronic structure of Li*<sub>7</sub>Cs (Cmmm) at 80 GPa e hibits interesting features of 1D-like bands. Its densit of states at the Cs $\delta 5s^{b}$  and Cs $\delta 5p^{b}$  bands [see Fig. 6(a)] displa s a remarkable feature of t pical 1D electronic structure, the emergence of

hich in a three-dimensional ervironment is rather une pected. Figure 6(b) sho s the band structure of Li<sub>7</sub>Cs (*Cmmm*) at 80 GPa. The Cs $\delta$ 5*s* $\flat$  and Cs $\delta$ 5*p* $\flat$  bands displa strong dispersion onl along the  $\Gamma$ -*X*<sub>5</sub> direction, but are nearl dispersionless along other directions. The 1D band structure is caused b the unique cr stal structure of Li<sub>7</sub>Cs (*Cmmm*) [see the inset of Fig. 3(a)]. The Cs chains are embedded in the three-dimensional Li net orks and are separated from each other. Along the chain direction, the Cs atoms are close to each other and their cores  $\sigma$ verlap.